Prospect for a ridge in p+Pb collisions

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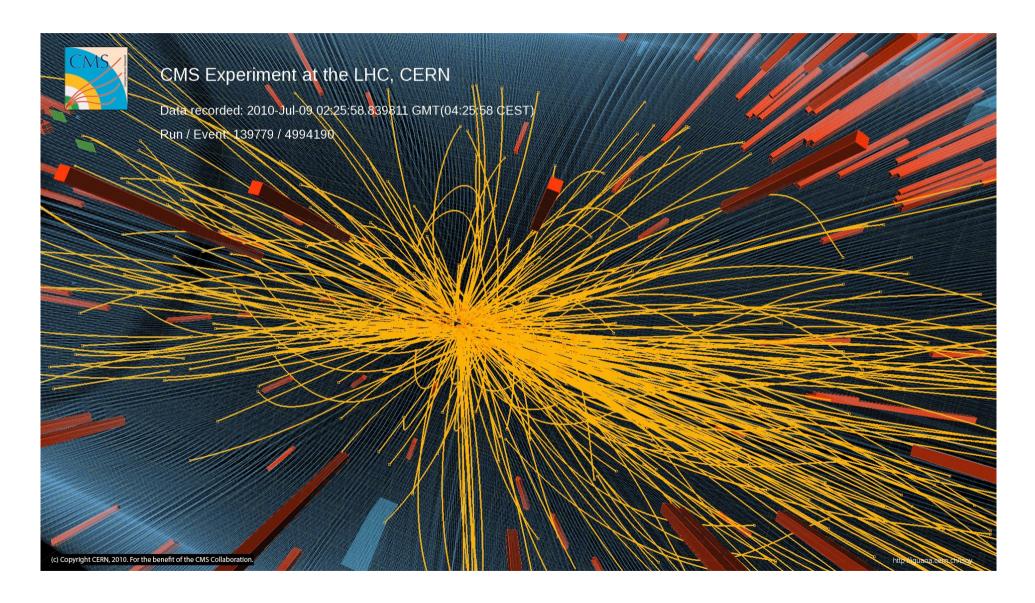
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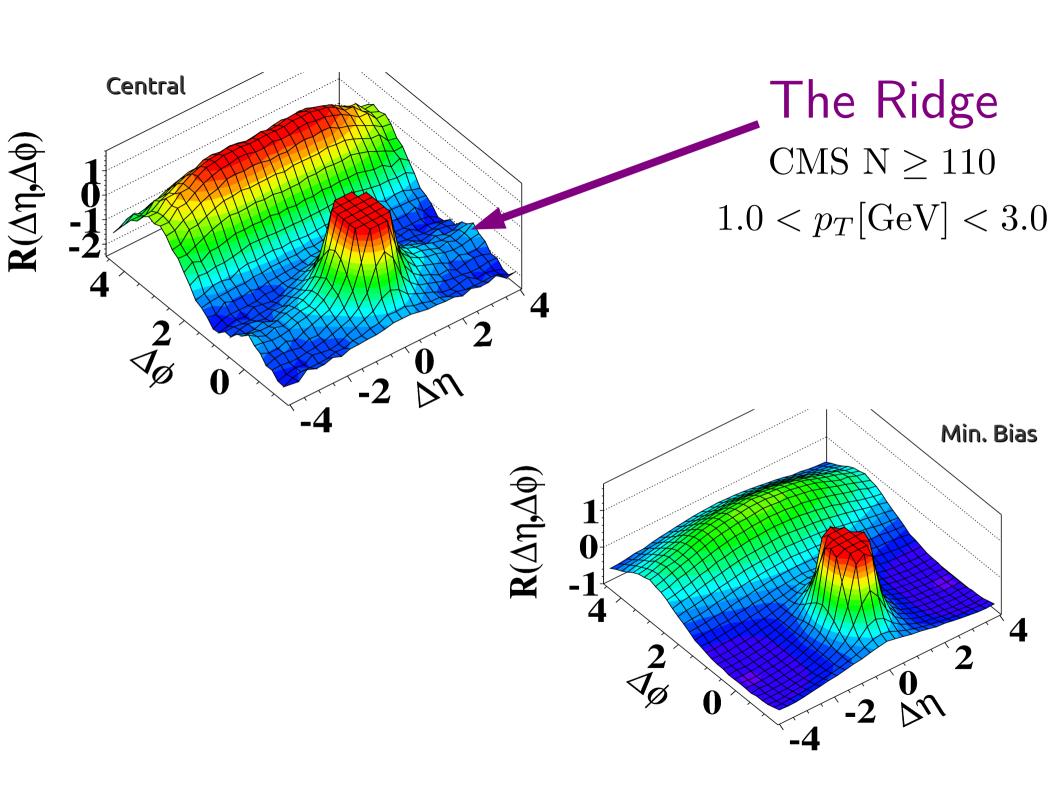
July 30th, 2012

NC STATE UNIVERSITY

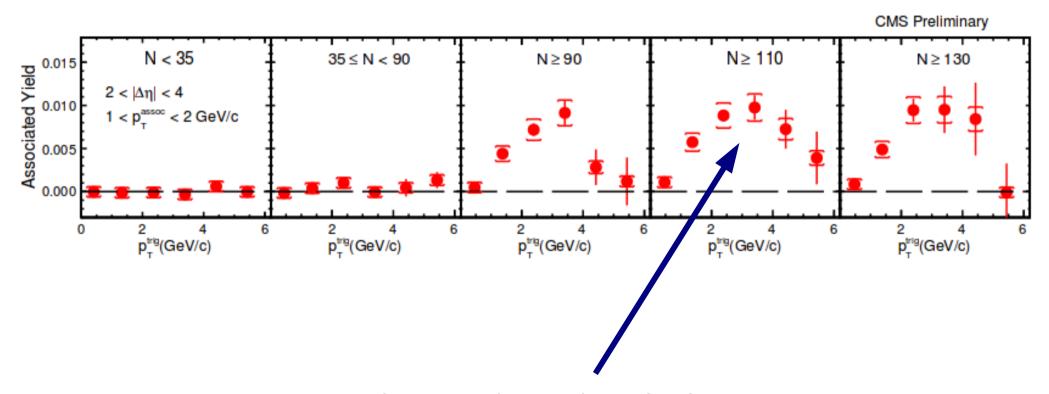
The Ridge



More than 200 charged particles!

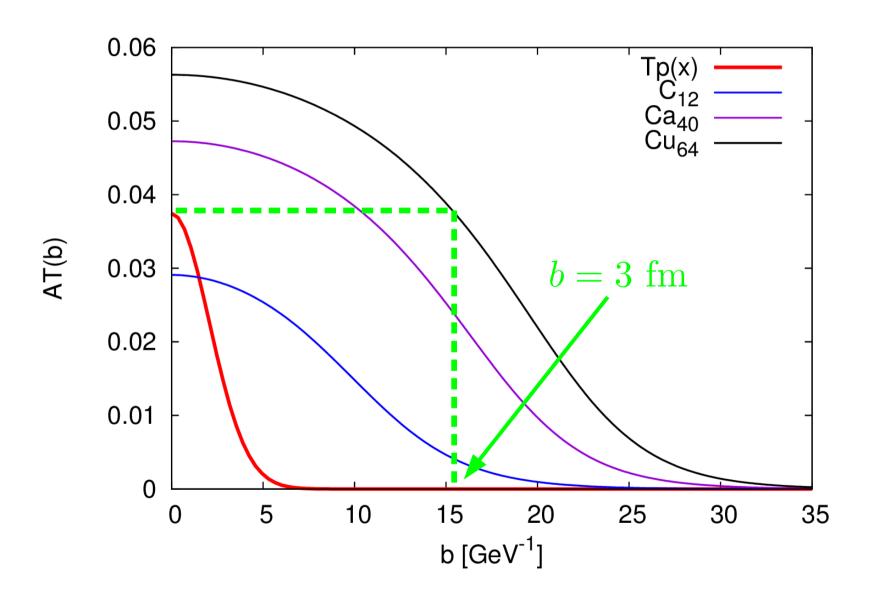


The Ridge

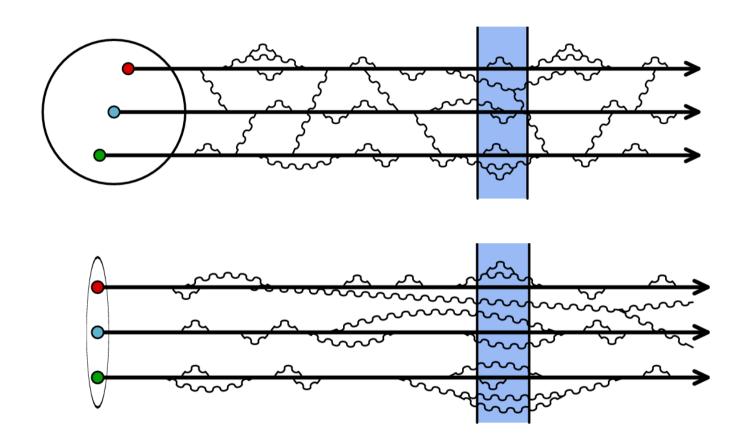


- 1. There is a clear scale in the data
- 2. It is semi-hard and will be argued to be Q_s

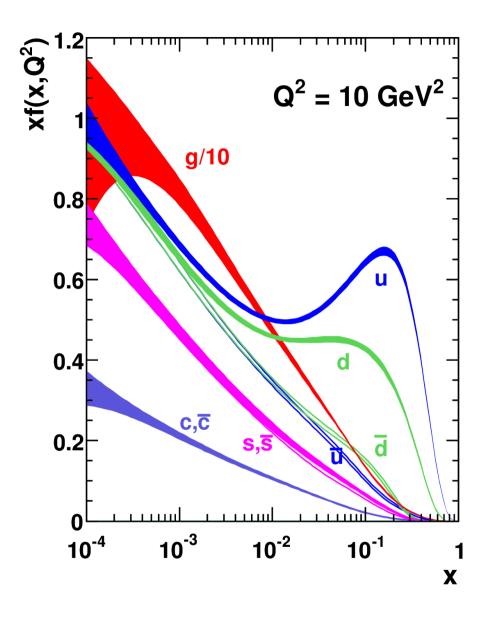
Multiplicity the same as in Cu+Cu!



Wave function of the proton



$$x \sim 10^{-2}$$
 at $\sqrt{s} = 200$ GeV
 $x \sim 10^{-4}$ at $\sqrt{s} = 7$ TeV



Growth of gluon distribution function at small x is seen experimentally.

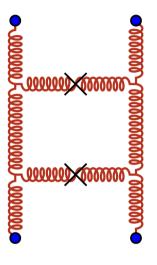
Data / figures from: http://mstwpdf.hepforge.org/

The role of STRONG color sources

Diagram:

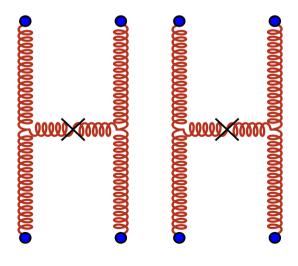
Min. Bias

Central



$$\mathcal{O}\left(\alpha_s^4\right)$$

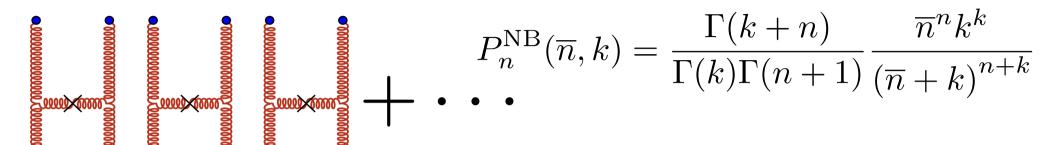
$$\mathcal{O}\left(1\right)$$



$$\mathcal{O}\left(\alpha_s^6\right)$$

$$\mathcal{O}\left(\alpha_s^{-2}\right)$$

High multiplicity are b=0 collisions



Dumitru, Gelis, McLerran, Venugopalan, NPA810 91-108 (2008). Dusling, Fernandez-Fraile, Venugopalan NPA828 (2009) 161-177. Gelis, Lappi, McLerran, NPA828 (2009) 149-160.

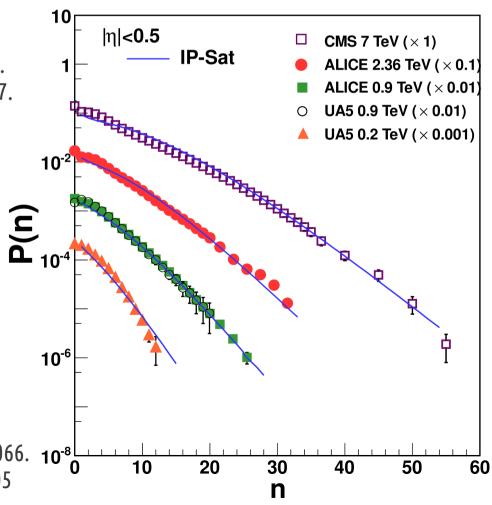
$$k = \zeta \frac{\left(N_c^2 - 1\right) S_\perp Q_s^2}{2\pi}$$

$$\zeta = 0.155$$
 [Empirical]

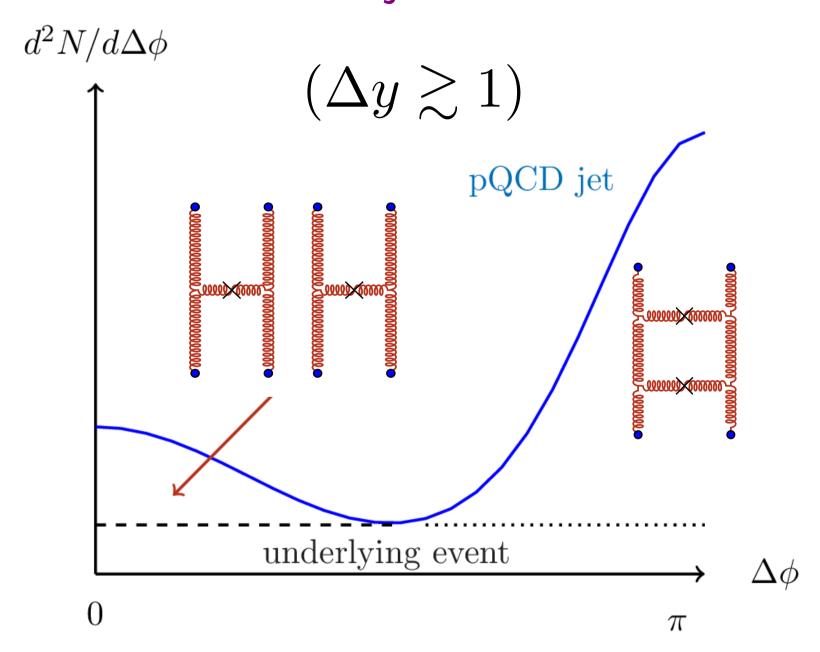
$$\zeta = 0.2 - 1.5$$
 [Lattice]

Emprical: Tribedy, Venugopalan, NPA850 (2011) 136-156.

Lattice (CYM): Lappi, Srednyak, Venugopalan, JHEP01 (2010) 066. 10-8L Schenke, Tribedy, Venugopalan, arXiv:1206.6805



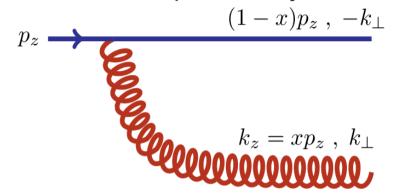
Forward jet structure



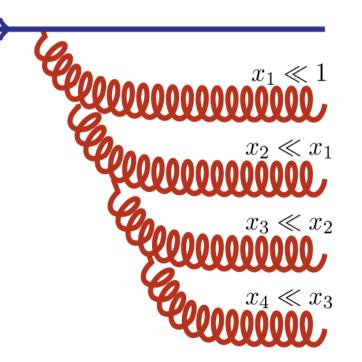
Gluon radiation

As the energy is increased new gluons are emitted with probability

$$\mathrm{d}P_{\mathrm{Brem}} \sim C_R \frac{\alpha_s}{\pi^2} \frac{\mathrm{d}^2 k_{\perp}}{k_{\perp}^2} \frac{\mathrm{d}x}{x}$$



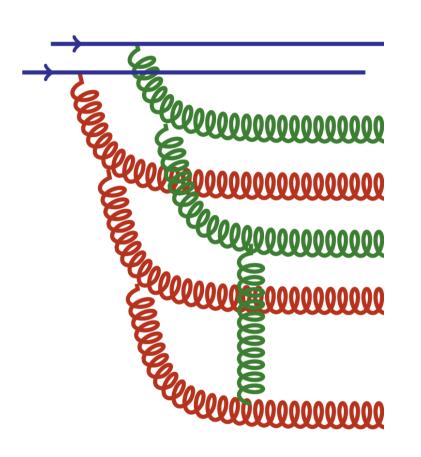
And as long as the density remains low the evolution is linear



$$\frac{\partial T(\mathbf{r}_{\perp}, Y)}{\partial Y} = \frac{\alpha_s N_c}{2\pi^2} \int d\mathbf{r}_1 \frac{\mathbf{r}_{\perp}^2}{\mathbf{r}_1^2 (\mathbf{r}_{\perp} - \mathbf{r}_1)^2} \times [T(\mathbf{r}_1, Y) + T(\mathbf{r}_{\perp} - \mathbf{r}_1, Y) - T(\mathbf{r}_{\perp}, Y)]$$

Kuraev, Lipatov, Fadin, Sov.Phys.JETP44 443-450 (1976). Sov.Phys.JETP45 199-204 (1977). Balitsky, Lipatov, Sov.J.Nucl.Phys 28 822-829 (1978).

BK Evolution Equation



Balitsky, NPB 463, 99 (1996). Kovchegov, PRD 60, 034008 (1999).

Jalilian-Marian, Kovner, McLerran, Weigert, PRD 55 5414 (1997).

Jalilian-Marian, Kovner, Leonidov, Weigert, NPB 504 415 (1997),

PRD 59 014014 (1999).

$$\frac{\partial T(\mathbf{r}_{\perp}, Y)}{\partial Y} = \frac{\alpha_s N_c}{2\pi^2} \int d\mathbf{r}_1 \frac{\mathbf{r}_{\perp}^2}{\mathbf{r}_1^2 (\mathbf{r}_{\perp} - \mathbf{r}_1)^2} \times [T(\mathbf{r}_1, Y) + T(\mathbf{r}_{\perp} - \mathbf{r}_1, Y) - T(\mathbf{r}_{\perp}, Y) - T(\mathbf{r}_1, Y)T(\mathbf{r}_{\perp} - \mathbf{r}_1, Y)]$$

NLO BK Equation

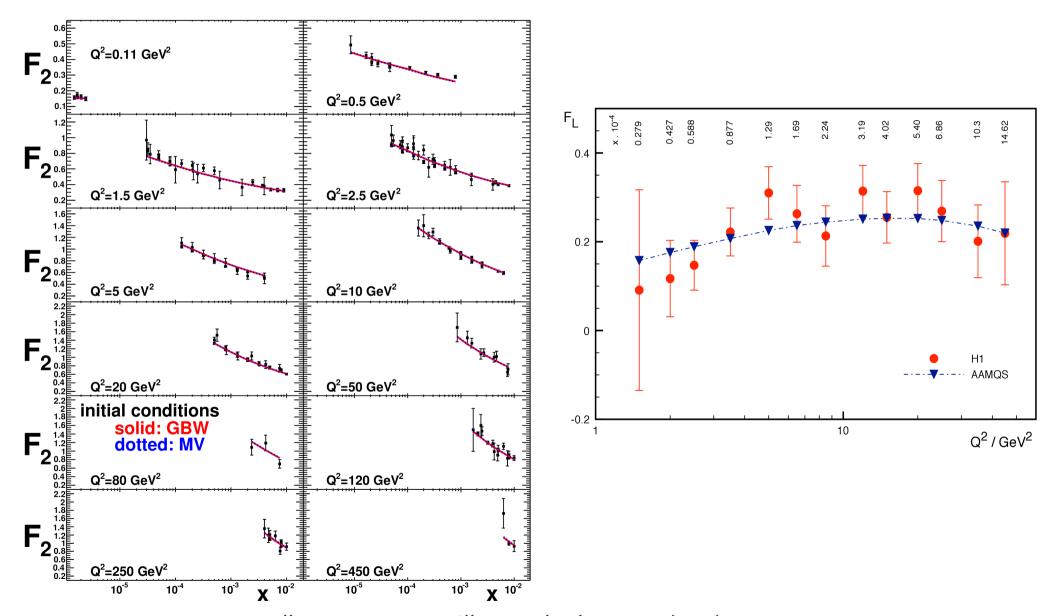
$$\frac{\partial T(\mathbf{r}_{\perp}, Y)}{\partial Y} = \int d\mathbf{r}_{1} \mathcal{K}_{\text{Bal.}}(\mathbf{r}_{\perp}, \mathbf{r}_{1}, \mathbf{r}_{\perp} - \mathbf{r}_{1}) \times$$

$$[T(\mathbf{r}_{1}, Y) + T(\mathbf{r}_{\perp} - \mathbf{r}_{1}, Y) - T(\mathbf{r}_{\perp}, Y) - T(\mathbf{r}_{1}, Y)T(\mathbf{r}_{\perp} - \mathbf{r}_{1}, Y)]$$

$$\mathcal{K}_{\text{Bal.}}(\mathbf{r},\mathbf{r}_1,\mathbf{r}_2) = \frac{\alpha_s(\mathbf{r})N_c}{\pi} \left[\frac{\mathbf{r}^2}{\mathbf{r}_1^2\mathbf{r}_2^2} + \frac{1}{\mathbf{r}_1^2} \left(\frac{\alpha_s(\mathbf{r}_1^2)}{\alpha_s(\mathbf{r}_2^2)} - 1 \right) + \frac{1}{\mathbf{r}_2^2} \left(\frac{\alpha_s(\mathbf{r}_2^2)}{\alpha_s(\mathbf{r}_1^2)} - 1 \right) \right]$$

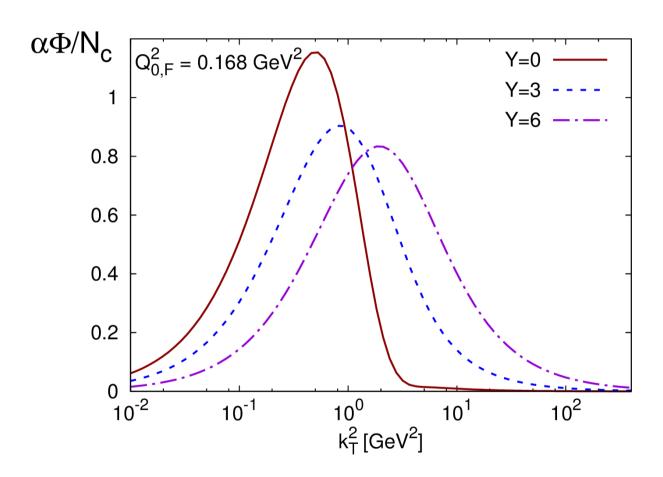
Balitsky, Chirilli PRD 77 014019 Kovchegov, Weigert NPA 784 188 Albacete, Kovchegov PRD 75 125021

Deep inelastic scattering on the Proton



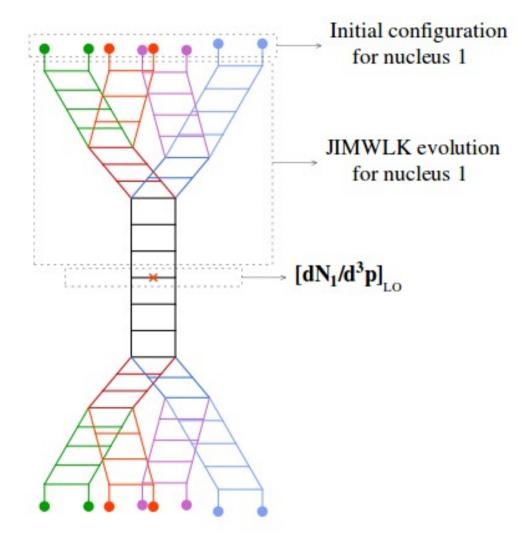
Albacete, Armesto, Milhano, Salgado; PRD80 (2009) 034031. Quiroga-Arias, Albacete, Armesto, Milhano, Salgado; PRD80 J.Phys.G G38 (2011) 124124.

Unintegrated gluon distribution



$$\Phi_{A_{1,2}}(x,k_{\perp}) = \frac{N_c k_{\perp}^2}{4 \alpha_s(\mathbf{k}_{\perp})} \int d^2 \mathbf{r}_{\perp} J_0(k_{\perp} r_{\perp}) \left[1 - T_{A_{1,2}}(r_{\perp}, \ln(1/x)) \right]^2$$

k_T factorization: single gluon production



$$\left\langle \frac{\mathrm{d}N_1}{\mathrm{d}^2 \mathbf{p}_{\perp} \mathrm{d}y_p} \right\rangle_{\mathrm{LLor}} = \frac{8\alpha_s(p_{\perp})S_{\perp}}{C_F(2\pi)^4} \frac{1}{\mathbf{p}_{\perp}^2} \int \frac{\mathrm{d}^2 k_{\perp}}{(2\pi)^2} \Phi_{A_1}(y_p, k_{\perp}) \Phi_{A_2}(y_p, p_{\perp} - k_{\perp})$$

k_T factorization: double gluon production

$$\left\langle \frac{\mathrm{d}N_{2}}{\mathrm{d}^{2}\mathbf{p}_{\perp}\mathrm{d}y_{p}\mathrm{d}^{2}\mathbf{q}_{\perp}\mathrm{d}y_{q}} \right\rangle_{\mathrm{LLog}} = \frac{32\alpha_{s}(\mathbf{p}_{\perp})\alpha_{s}(\mathbf{q}_{\perp})}{(2\pi)^{10}N_{c}C_{F}^{3}\zeta} \frac{1}{\mathbf{p}_{\perp}^{2}\mathbf{q}_{\perp}^{2}}$$

$$\times \left\{ \int \mathrm{d}^{2}\mathbf{k}_{1\perp}\Phi_{A_{1}}^{2}(y_{p},\mathbf{k}_{1\perp})\Phi_{A_{2}}(y_{p},\mathbf{p}_{\perp}-\mathbf{k}_{1\perp})\Phi_{A_{2}}(y_{q},\mathbf{q}_{\perp}+\mathbf{k}_{1\perp})$$

$$+ \int \mathrm{d}^{2}\mathbf{k}_{1\perp}\Phi_{A_{1}}^{2}(y_{p},\mathbf{k}_{1\perp})\Phi_{A_{2}}(y_{p},\mathbf{p}_{\perp}-\mathbf{k}_{1\perp})\Phi_{A_{2}}(y_{q},\mathbf{q}_{\perp}-\mathbf{k}_{1\perp})$$

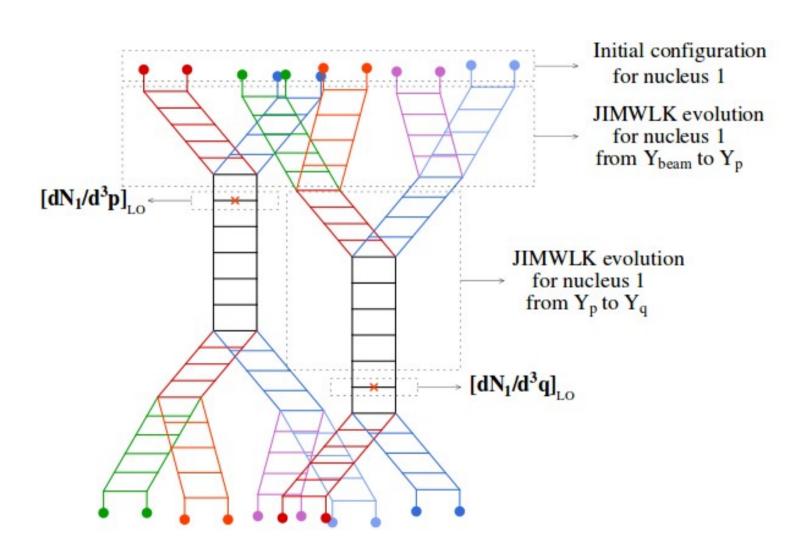
$$+ \int \mathrm{d}^{2}\mathbf{k}_{1\perp}\Phi_{A_{2}}^{2}(y_{q},\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{p},\mathbf{p}_{\perp}-\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{q},\mathbf{q}_{\perp}+\mathbf{k}_{1\perp})$$

$$+ \int \mathrm{d}^{2}\mathbf{k}_{1\perp}\Phi_{A_{2}}^{2}(y_{q},\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{p},\mathbf{p}_{\perp}-\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{q},\mathbf{q}_{\perp}-\mathbf{k}_{1\perp})$$

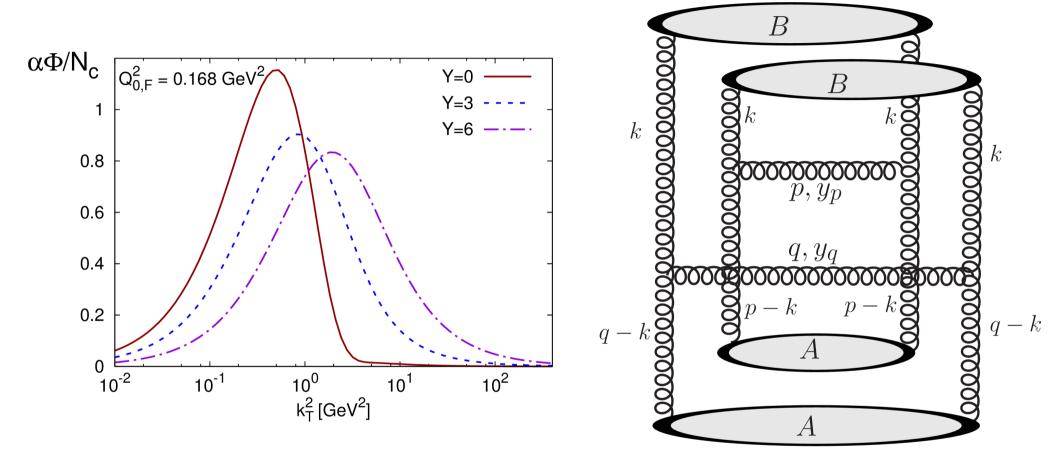
$$+ \int \mathrm{d}^{2}\mathbf{k}_{1\perp}\Phi_{A_{2}}^{2}(y_{q},\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{p},\mathbf{p}_{\perp}-\mathbf{k}_{1\perp})\Phi_{A_{1}}(y_{q},\mathbf{q}_{\perp}-\mathbf{k}_{1\perp})$$

Gelis, Lappi, Venugopalan, PRD78, 050419 (2008).
PRD78, 054020 (2008).
PRD79, 094017 (2009).
Dusling, Gelis, Lappi, Venugopalan, NPA 836 159-182 (2010).

k_T factorization: double gluon production



Angular Structure

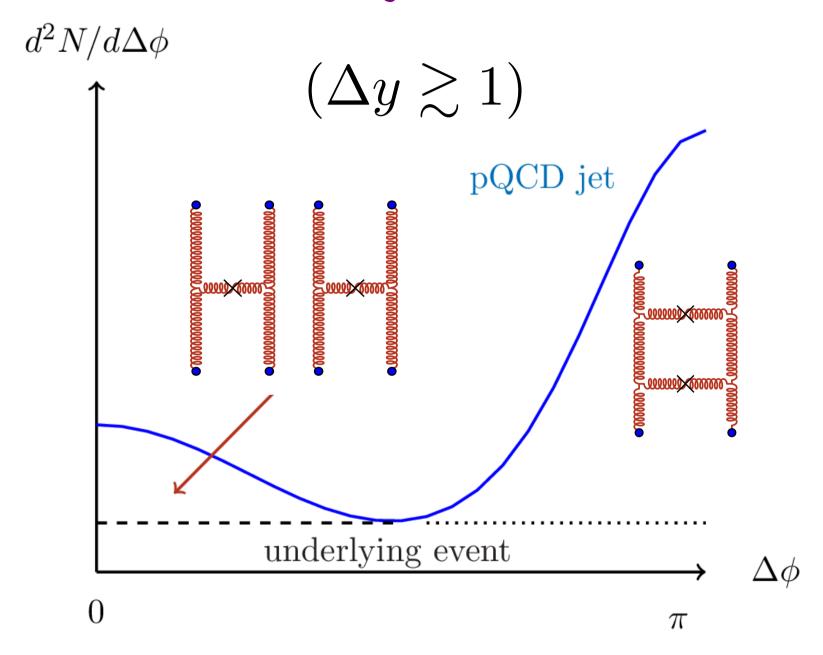


Condition for Ridge (Qualitatively):

$$|\mathbf{k}_{\perp}| \sim |\mathbf{p}_{\perp} - \mathbf{k}_{\perp}| \sim |\mathbf{q}_{\perp} \pm \mathbf{k}_{\perp}| \sim Q_s$$

Dumitru, Dusling, Gelis, Lalilian-Marion, Lappi, Venugopalan, PLB 697 12-25 (2011).

Forward jet structure

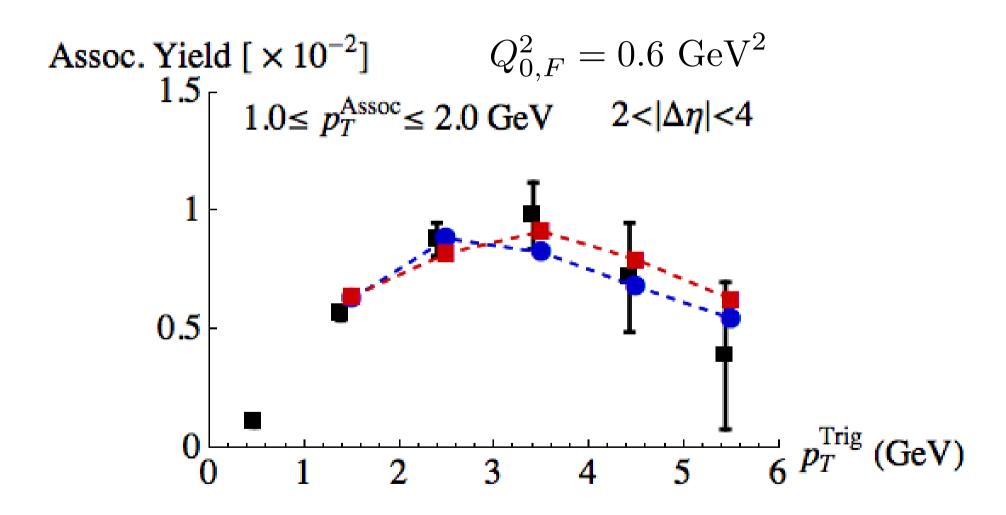


Centrality Dependence

Assoc. Yield [$\times 10^{-2}$]

$$Q_{0,F}^2(x_0 = 0.01) = 0.15, 0.3, 0.45, 0.6 \text{ GeV}^2$$

Trigger Dependence



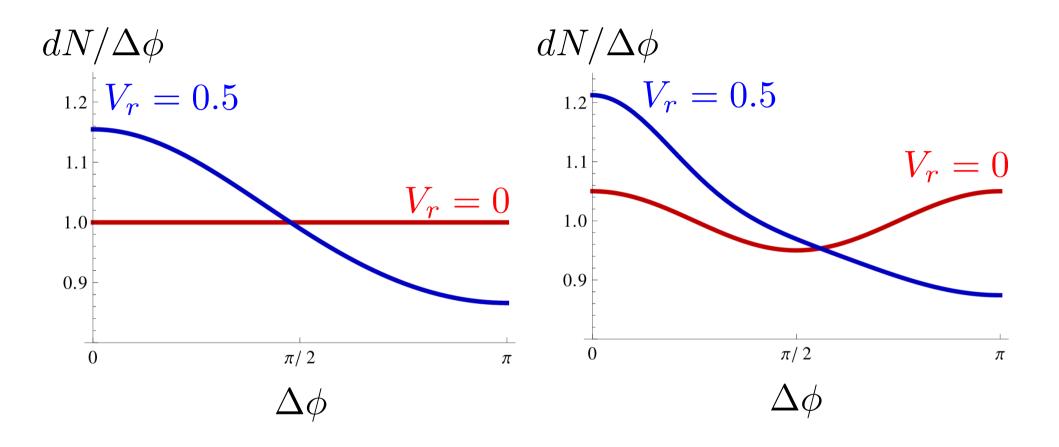
Blast Wave I

$$\frac{d^{2}N}{d\Delta\phi} = \int_{-\pi}^{\pi} d\Psi \mathcal{J}\left(\Psi, \Delta\phi\right) \frac{d^{2}N}{d\Delta\tilde{\phi}} \left(\Delta\tilde{\phi}\left(\Psi, \Delta\phi\right)\right)$$

$$\begin{split} 2\sin^2\left(\frac{\Delta\tilde{\phi}}{2}\right) &= \\ \frac{\sqrt{1-\beta^2}\left(1-\cos\left(\Delta\phi\right)\right)}{1-2\beta\cos\Psi\cos\left(\frac{\Delta\phi}{2}\right)+\frac{\beta^2}{2}\left(\cos\left(\Delta\phi\right)+\cos\left(2\Psi\right)\right)} \;. \end{split}$$

$$\mathcal{J} = \frac{1 - \beta^2}{\left(1 - \beta\cos\left(\Psi + \Delta\phi/2\right)\right)\left(1 - \beta\cos\left(\Psi - \Delta\phi/2\right)\right)},$$

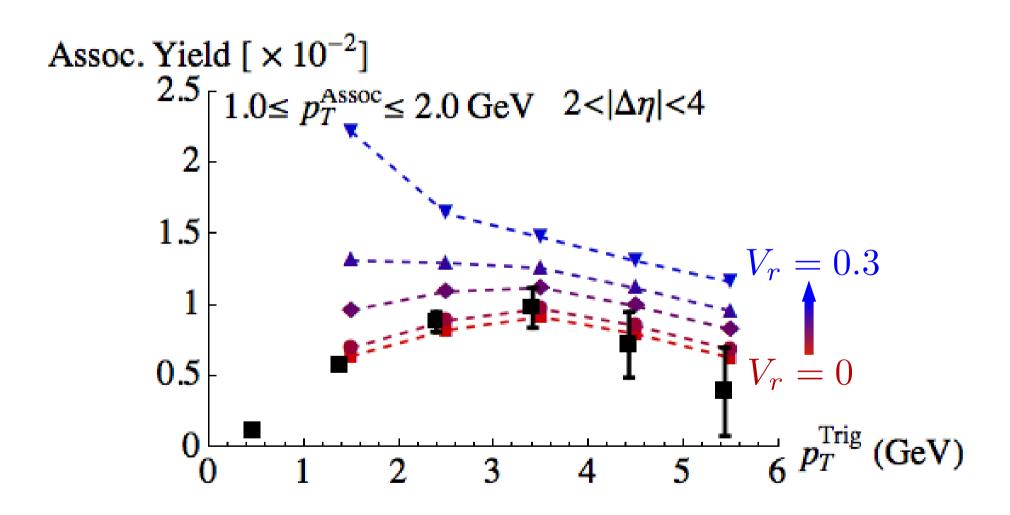
Blast Wave II



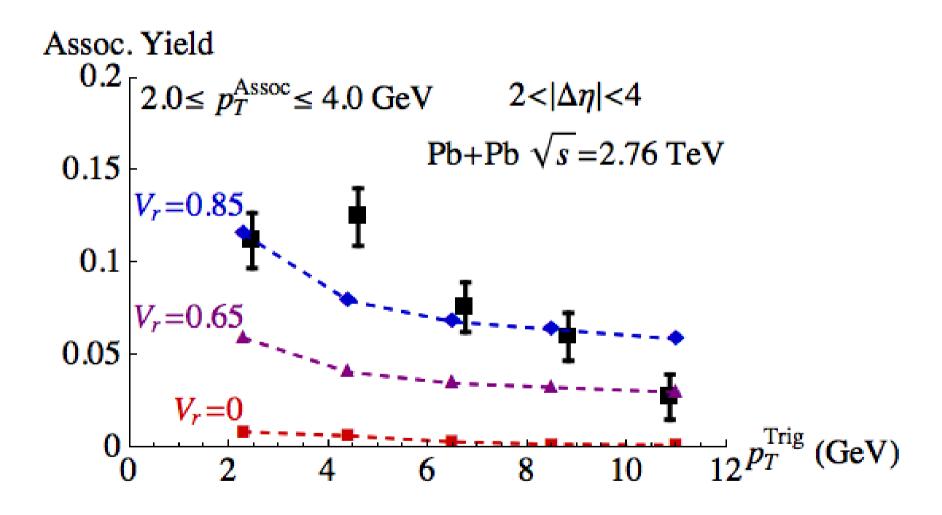
Left: No intrinsic correlation in $\Delta \phi$ followed by radial boost.

Right: Intrinsic azimuthal correlation followed by boost.

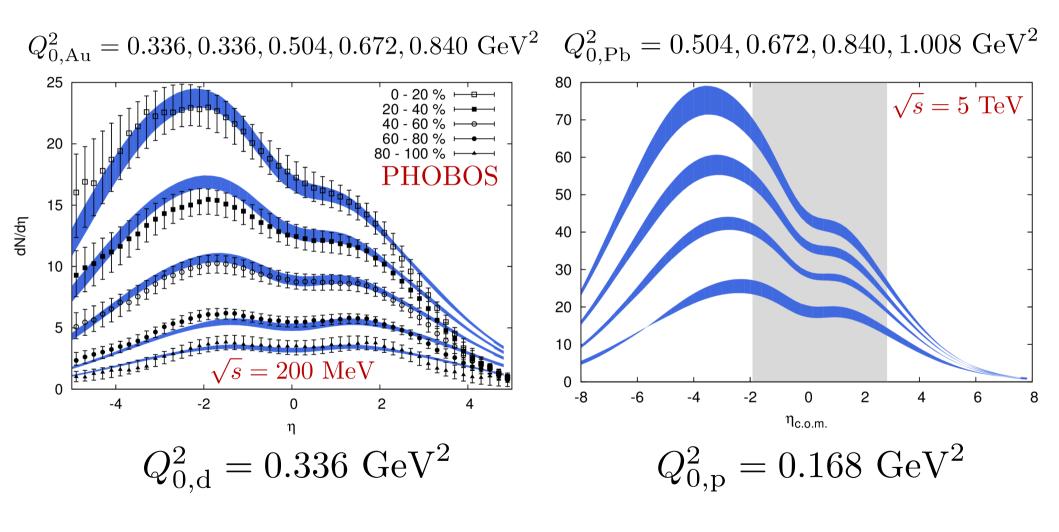
Blast wave results



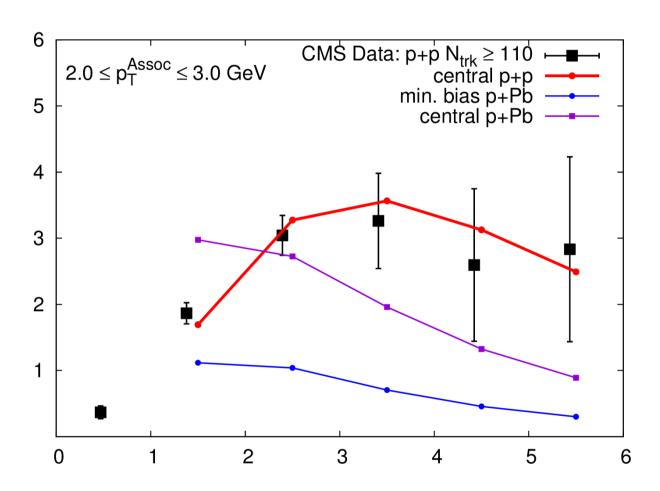
Results for Pb+Pb



Multiplicity in p+Pb at 5 TeV

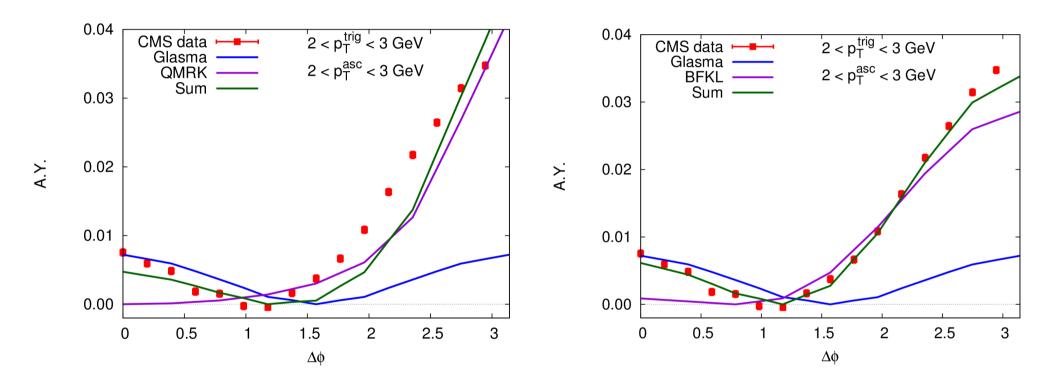


Ridge in p+Pb



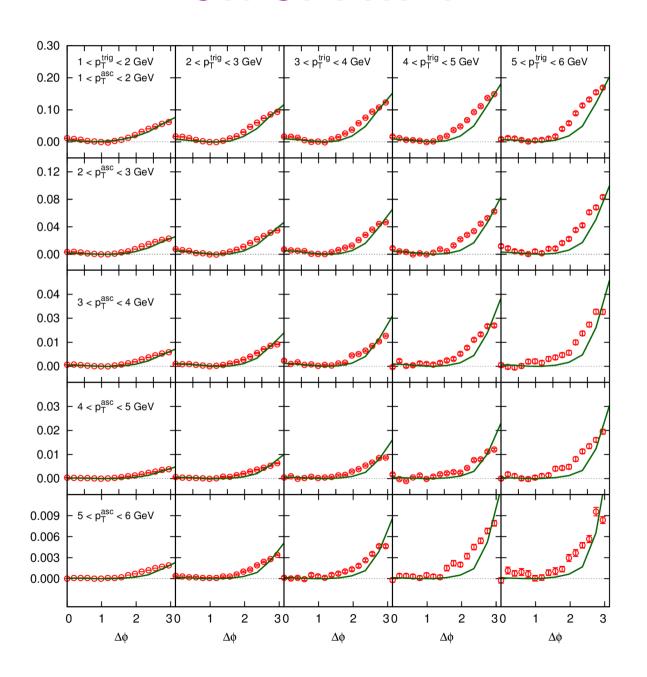
Ridge in p+Pb is smaller than in p+p for CMS acceptance. Signal will also have to be pulled from a larger background.

Understanding the away-side



There is a clear need for evolution between the triggered particles (even for a rapidity gap as small as 2-4 units)

Jet Structure



Summary

• Strong color sources lead to α_s^8 enhancement of QCD diagram responsible for near-side enhancement

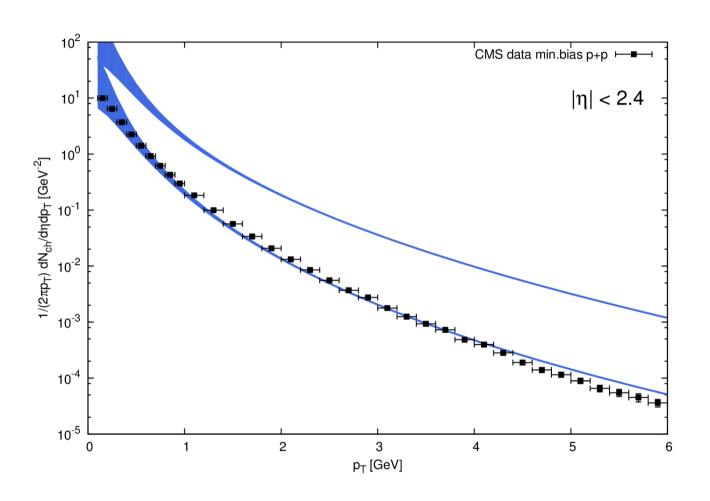
Structure of ridge correlation constrains radial flow in p+p

Radial flow explains identical measurements in Pb+Pb

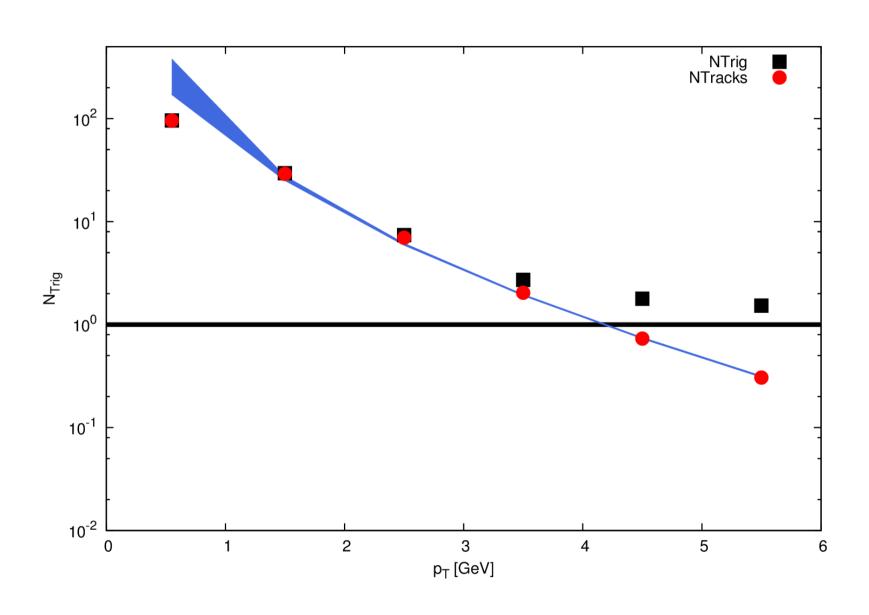
Ridge tougher to see in asymmetric collisions

Backup

p+p p_T distribution



CMS Acceptance



BFKL Formalism

$$\begin{split} \frac{d^{2}N_{AB}}{d^{2}\mathbf{p}_{\perp}d^{2}\mathbf{q}_{\perp}dy_{p}dy_{q}} &= \frac{32N_{c}\alpha_{s}(\mathbf{p}_{\perp})\alpha_{s}(\mathbf{q}_{\perp})}{(2\pi)^{8}C_{F}} \frac{1}{\mathbf{p}_{\perp}^{2}\mathbf{q}_{\perp}^{2}} \\ &\times \int d^{2}\mathbf{k}_{0\perp} \int d^{2}\mathbf{k}_{3\perp}\Phi_{A}(x_{1},\mathbf{k}_{0\perp})\Phi_{B}(x_{2},\mathbf{k}_{3\perp})\mathcal{G}(\mathbf{k}_{0\perp}-\mathbf{p}_{\perp},\mathbf{k}_{3\perp}+\mathbf{q}_{\perp},y_{p}-y_{q}) \\ \mathcal{G}(a,b,\Delta y) &= \frac{1}{(2\pi)^{2}} \frac{1}{(a^{2}b^{2})^{1/2}} \sum_{n} e^{in\overline{\phi}} \int_{-\infty}^{+\infty} d\nu \ e^{\omega(\nu,n)\Delta y} e^{i\nu\ln(a^{2}/b^{2})} \\ \omega(\nu,n) &= -2\overline{\alpha}_{s} \operatorname{Re} \left[\Psi\left(\frac{|n|+1}{2}+i\nu\right) - \Psi(1) \right] \\ \overline{\alpha}_{s} &\equiv \frac{N_{c}\alpha_{s}\left(\sqrt{ab}\right)}{\pi} \\ \overline{\phi} &\equiv \arccos\left(\frac{a\cdot b}{|a| |b|}\right) \end{split}$$